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# Light-controlled microwave whispering-gallery-mode quasi-optical resonators at 50W LED array illumination

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We present experimental observations of light-controlled resonance effects in microwave whispering-gallery-mode quasi-optical dielectric-semiconductor disk resonators in the frequency band of 5 GHz to 20 GHz arising due to illumination from a light emitting diode (LED) of 50W power range. We obtain huge enhancement of photo-sensitivity (growing with the resonator  $Q$ -factor) that makes light-microwave interaction observable with an ordinary light (no laser) at conventional brightness (like an office lighting) in quasi-optical microwave structures at rather long (centimeter-scale) wavelength. We also demonstrate non-conventional photo-response of Fano resonances when the light suppresses one group of resonances and enhances another group. The effects could be used for the optical control and quasi-optical switching of microwave propagation through either one or another frequency channel. © 2015 Author(s). All article content, except where otherwise noted, is licensed under a Creative Commons Attribution 3.0 Unported License. [<http://dx.doi.org/10.1063/1.4928868>]

## I. INTRODUCTION

Whispering-gallery-mode (WGM) resonators found numerous applications<sup>1</sup> due to their exceptional performance associated with extremely high quality factors  $Q$ . Major progress is achieved in optics where  $Q$  factors exceeding one billion<sup>2</sup> are demonstrated and huge enhancement to other characteristics of devices are realized. The devices include biological and chemical sensors of nanoparticles and bio-molecules,<sup>3,4</sup> unique ultralow-threshold lasers,<sup>5</sup> advanced opto-electronic and electro-optical switchers and oscillators,<sup>6–8</sup> etc.

Microwave WGM structures do also have significant  $Q$  factors ( $Q > 10^5$  at the room temperature<sup>9</sup>) that allows for realization of high-performance multi-pole filters,<sup>10</sup> enhanced microwave oscillators,<sup>11</sup> ultra-sensitive (nanoWatt level) photonic microwave receivers,<sup>12</sup> new dielectric material characterization tools and techniques.<sup>13,14</sup>

In this paper, we present experimental observations of light-controlled resonance effects in microwave whispering-gallery-mode (WGM) disk resonators in the frequency band of 5 GHz to 20 GHz arising due to illumination from a light emitting diode (LED) integrated circuit array of 50W power range.

Light-controlled microwave devices offer the advantages of very high bandwidth and unique functionality, particularly, in quasi-optical systems in the higher frequency bands.<sup>15–17</sup> Optically controlled dielectric WGM resonators have also been considered in this respect.<sup>18</sup> In distinction from numerical simulations in Ref. 18 where signal transmission peaks are of interest (the resonator works as a pass-band filter), we present experimental study of those cases when the resonant deeps in signal transmission are the main feature (the structure is used as a stop-band filter device).

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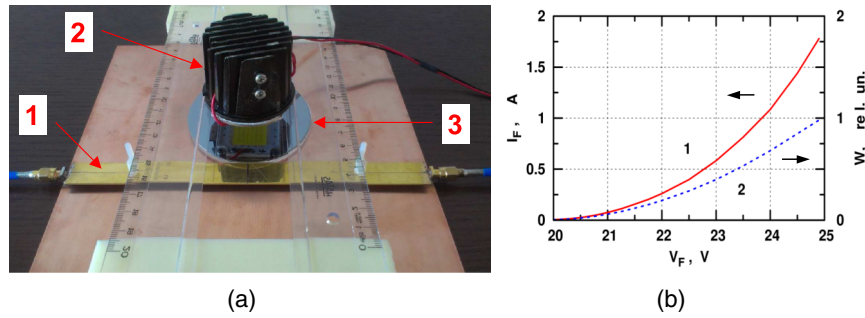


FIG. 1. (a) Experimental setup including (1) strip line, (2) LED array with a heat sink, and (3) WGM resonator, and (b) LED array current-voltage (curve 1) and luminosity-voltage (curve 2) characteristics.

## II. EXPERIMENTAL SETUP

A general view of experimental setup is shown in Fig. 1. The system consists of a cylindrical disk resonator placed on a metallic plate, a suspended strip line for the microwave excitation, and a 50W LED array (P003F13W-50B6C8) for the illumination of resonator surface. Metallic plate works as a mirror supporting the resonances at the lower frequencies as those in the disk suspended in free space (yet, they are higher as compared to those in the disk squeezed between two plates<sup>19</sup>). Microwave measurements are made with the E5071C network analyzer in the frequency band from 300 KHz to 20 GHz. The measured quantity is the S21 coefficient representing the signal transmission through the system as shown in Fig. 2.

Each resonator is assembled as a stack of disk wafers of sapphire ‘s’, fused quartz ‘q’, and high-resistivity silicon ‘si’ ( $\rho = 30 \text{ KOhm cm}$ ) which is, typically, used as a top layer. The stacks are denoted as, e.g.,  $MsNsi$  and  $MqNsi$  where the coefficients  $M$  and  $N$  specify the numbers of relevant wafers. The thickness of ‘s’, ‘q’, and ‘si’ wafers is  $t_s = 0.40 - 0.42 \text{ mm}$ ,  $t_q = 0.52 \text{ mm}$ , and  $t_{si} = 0.50 \text{ mm}$ , respectively, and the diameters are  $D_s = D_{si} = 76.2 \text{ mm}$  and  $D_q = 75.0 \text{ mm}$ . Below we present the results for  $7s1si$ ,  $7q3si$ ,  $10s1si$ , and  $7q2si$  stacks.

The size and the number of wafers in each resonator have been chosen so as to provide a well-developed resonant spectrum in the given frequency band. The bulk of the resonator consists of traditional low-loss materials whereas semiconductor wafers are used only as the top layers to be illuminated. Both the depths and the frequencies of the resonant deeps could be fine tuned by adjusting the number of wafers and the size of the air slot between the resonator and the excitation line. The metallic plate used as a substrate helps to provide an efficient excitation of the resonator. In addition, once the plate is operating as a mirror, when the top semiconductor wafer is illuminated, the image wafer also increases the conductivity and contributes to the better control of the resonator performance.

## III. DARK-STATE WGM RESONATOR RESPONSE

Fig. 3 shows the dark-state microwave transmission spectra of strip line loaded with one of two kinds of disk resonators,  $7q3si$  and  $7s1si$ . Transmission of the strip line with no resonator is

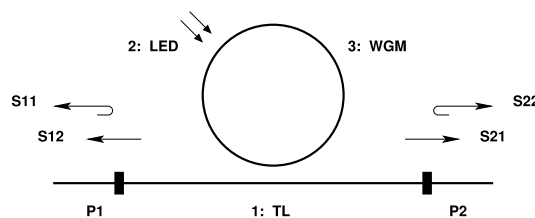


FIG. 2. Schematics of the system with (1) transmission line (TL), (2) LED array, and (3) WGM resonator.

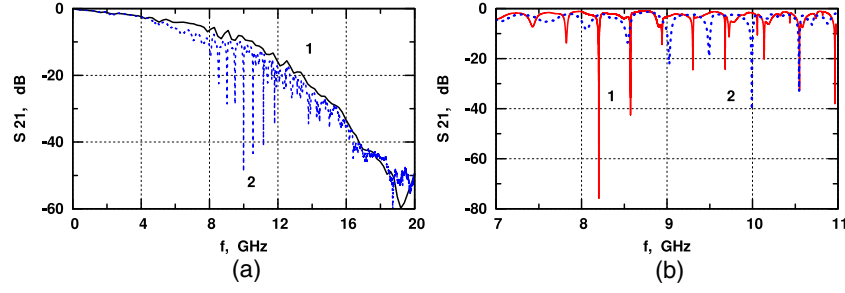


FIG. 3. Dark-state microwave transmission spectra  $S_{21}$  in case of (a) no resonator (curve 1, to be used as a reference) and  $7q3si$  resonator (curve 2) and (b)  $7s1si$  and  $7q3si$  resonators when presented with respect to the reference (curves 1 and 2, respectively).

shown by curve 1 in Fig. 3(a). The line is made to have the intrinsic impedance of 50 Ohms at the frequency  $f = 10$  GHz. Transmission of the line decreases with increasing the frequency due to radiation losses, dropping down to -60 dB near the frequency of  $f = 20$  GHz.

When loaded with a disk resonator, the strip line shows the microwave transmission spectrum with a sequence of resonant deeps due to the WGM disk excitation (curve 2 in Fig. 3(a)). Further, we show the  $S_{21}$  spectra with respect to the reference curve of unloaded line so as to better represent the resonator effects.

Exact simulations of three-dimensional WGM disk resonators are complicated,<sup>18,20,21</sup> especially, with account of excitation lines. In the meantime, simplistic estimates based on two-dimensional model<sup>22</sup> confirm the excitation of basic modes at about those frequencies as observed in the experiments. Fine tuning of disk position with respect to strip line, while keeping a minor slot between the disk and the line, allows us to obtain a significant depth of the resonant deeps in the signal transmission.

In our cases of  $7s1si$  and  $7q3si$  resonators (curves 1 and 2 in Fig. 3(b), respectively), fine tuning resulted in the leading resonant deeps of  $S_{21}$  signal down to the level of  $P_1 = -76$  dB and  $P_2 = -42$  dB with respect to the reference curve at the frequencies of  $f_1 = 8.20095$  GHz and  $f_2 = 9.99214$  GHz where the deeps are specified by the quality factors  $Q_1 = 3 \times 10^5$  and  $Q_2 = 5 \times 10^4$ , respectively. The deep quality factor  $Q$  is defined as  $Q = f/\Delta f$  where  $\Delta f$  is the width of the deep measured at the level of 3 dB above the minimum (measurements are made under the conditions of sufficient darkness).

#### IV. LIGHT SENSITIVITY OF TOP-SEMICONDUCTOR DISK RESONATOR

Light effects on the disk resonators with top semiconductor wafers (Figs. 4–6) were investigated using a white-color LED array whose current-voltage ( $I_F$ - $V_F$ ) and luminosity-voltage ( $W$ - $V_F$ ) characteristics are shown in Fig. 1(b). The characteristics were measured in a pulse mode with the

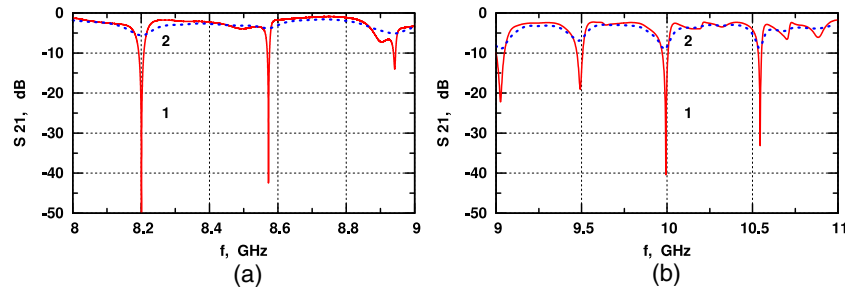


FIG. 4. Light sensitivity of resonances in case of (a)  $7s1si$  and (b)  $7q3si$  disk resonators (curves 1 and 2 show the dark and light states,  $V_F = 0$  V and  $V_F = 28$  V, respectively).

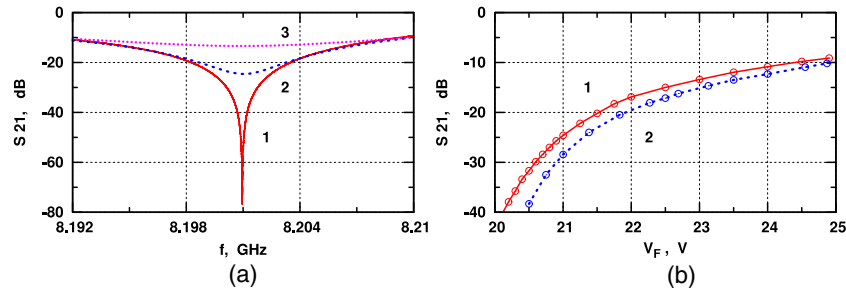


FIG. 5. (a) Light sensitivity of the leading resonant deep of S21 signal in case of  $7s1si$  resonator ( $V_F = 0, 21, 23V$ , curves 1 to 3, respectively) and (b) the LED voltage dependence of the leading resonant deeps in case of  $7s1si$  and  $7q3si$  structures.

light pulse duration  $\tau_p = 1s$  being the same as needed for recording separate traces of microwave spectrum. Microwave measurements were also made in a pulse mode, with sufficient relaxation time between pulses for avoiding heat and light accumulation effects, and with proper averaging of recorded signal over multiple traces.

Fig. 4 shows the light effects on the WGM resonances in case of  $7s1si$  and  $7q3si$  structures at the maximum power of LED array. In this case, the light suppresses all the resonances. This happens due to photo-excitation of electron-hole plasma in a semiconductor that induces electrical conductivity and creates significant microwave losses so as the resonances are completely damped.

Fig. 5 shows the effect of intermediate light intensity on the leading resonance of  $7s1si$  resonator and the LED voltage dependence of the leading resonant deeps in case of  $7s1si$  and  $7q3si$  structures ( $f_1 = 8.20095$  GHz and  $f_2 = 9.99214$  GHz, respectively). As expected, the resonance of higher quality factor  $Q$  and deeper drop of S21 signal is more sensitive to the illumination.

As an example, the illumination by the LED array at the voltage of  $V_F = 21V$  increases the microwave transmission at the leading resonant frequencies by 50 dB in case of  $7s1si$  resonator (curve 2 in Fig. 1(a)) and only by 13 dB in case of  $7q3si$  structure. The illuminance  $E_v$  on a silicon wafer (18 mm away from the array with active area of  $15 \times 15$  mm<sup>2</sup>) is estimated to be about 65 lx when using the LED array specification (the luminous flux is about  $F_v = 2370$  lm at  $V_F = 25$  V and  $I_F = 1.75$  A, that means one relative unit of luminosity  $W$  in Fig. 1(b)). Direct measurements with a pin photodiode SFH 203P show that  $E_v$  on the wafer at  $V_F = 21V$  is about 1.3 times greater than the illuminance in an office with a typical artificial lighting.

Finally, Fig. 6 shows the light effects in two other cases of resonators,  $10s1si$  and  $7q2si$  ones. Each resonator is tuned to have a weak coupling to the excitation line that results in a regular sequence of resonant deeps in microwave transmission as a function of frequency.

It is worth noting that, in the resonators with sapphire wafers, we observed a specific long-term photorefractive effect which did not appear in other cases. On a time-scale of a few seconds after the turning on the light (faster at the higher light intensity), the resonant deeps shifted towards the higher frequencies and showed increased microwave losses as compared to the values emerging at

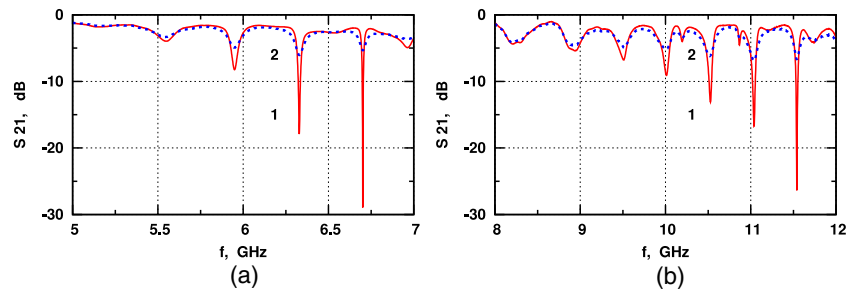


FIG. 6. Microwave transmission spectra in case of weak coupling of (a)  $10s1si$  and (b)  $7q2si$  resonators (curves 1 and 2 correspond to the dark and light states,  $V_F = 0V$  and  $V_F = 28V$ , respectively).

the start of the light pulse. Then, after the switching off the light, the relaxation to the initial state required a few seconds or, at the higher light intensity, up to a few minutes (both the rise and fall times of light pulses were about 30 ns).

The effect means a photo-induced reduction of sapphire dielectric constant that may be caused by photo-ionization of point-defect centers in the sapphire lattice. It also shows the accumulation tendency in a long sequence of light pulses. So, for avoiding the effect in our measurements, we used a low-duty pulse mode as explained earlier.

Enhanced light sensitivity of resonant deeps in microwave transmission provides a possibility of light-controlled switching of signal propagation with quasi-optical WGM resonators. At present, we could achieve switching with a 50W LED array at a partial load when the radiant flux is about  $E = 0.5\text{W/cm}^2$  (that corresponds to the illuminance of  $E_v = 65\text{ lx}$  at the given LED spectrum), though it is limited in the repetition rate  $f_R$  by the recombination lifetime in a silicon wafer (the latter is about  $\tau_R \sim 30\mu\text{s}$ , so as  $f_R$  should be less than  $\tau_R^{-1} \sim 30\text{kHz}$ ).

When using a semiconductor with recombination lifetime  $\tau_R \sim 1\text{ns}$ , the light-controlled switching could be obtained at the higher repetition rate with pulses of short duration  $\tau_p \sim 1\text{ns}$  at the pulse power flux of  $15\text{kW/cm}^2$ . In this case, the total radiant energy of a single pulse would be about  $7 \times 10^{-4}\text{J}$  that could be achieved with a short-pulse laser operating in a pulse mode at the repetition rate being about  $f_R \sim 1\text{GHz}$ .

## V. OVER-COUPLING EFFECTS

A group of more unusual effects is observed under the strong coupling of resonators with an excitation line. Fig. 7 shows the effects arising due to the strong coupling when using the same resonators as those in Fig. 6. In this case, when reducing the air slot between the disks and the strip line, the resonant lines acquire specific asymmetric shapes characteristic for the Fano resonances.<sup>23</sup> Fano resonances are the form of resonant oscillations emerging in the presence of continuous wave propagation when there is a coupling of oscillations with propagating waves.

Once the excitation line is supporting the propagating waves coupled to the resonator, the resonances being observed are always the Fano resonances which, though, look like conventional ones in case of a weak coupling. With increasing the coupling, they reveal their features as Fano resonances, such as the asymmetric shape of the resonant line profile.

When the coupling is particularly strong, the spectrum alters drastically and abruptly. This is the case of over-coupling, which is realized when the disks may touch the strip line, or slip under the strip, or one disk may be shifted with respect to another and partially overlay the strip line, etc. In these cases, the spectrum is getting especially complicated.

Regarding the light effects, the structures with over-coupling could be used as the light-controlled switches between different propagation channels. Fig. 7 shows a possibility of switching between two frequency channels through opening one channel and blocking the other (either  $f_1$  or  $f_2$ ) when turning on and off the light incident on the resonator ( $f_1 = 5.56\text{ GHz}$ ,  $f_2 = 6.64\text{ GHz}$  in case of  $10s1si$  structure and  $f_1 = 9.82\text{ GHz}$ ,  $f_2 = 9.52\text{ GHz}$  in the  $7q2si$  resonator).

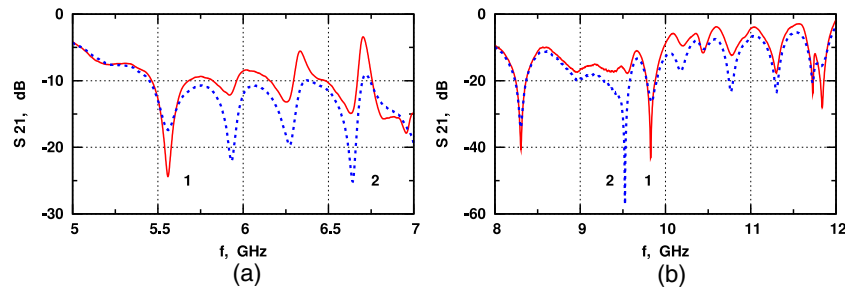


FIG. 7. Over-coupling effects in case of (a)  $10s1si$  and (b)  $7q2si$  resonators (curves 1 and 2 correspond to the dark and light states,  $V_F = 0\text{V}$  and  $V_F = 28\text{V}$ , respectively).

## VI. CONCLUSION

We present experimental study of light-controlled resonance response of microwave structures based on sapphire and fused quartz disk resonators with top semiconductor (intrinsic silicon) wafers being illuminated by a 50W LED array and operated in the frequency band of 5–20 GHz. Quantitative characteristics of the effect have been measured in the system configuration with resonant dips in the transmission signal. The effect could be used for light-controlled switching of signal propagation in this kind of microwave systems.

## ACKNOWLEDGMENTS

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